Numerical Solution to the Growth of Spherical Precipitates with Capillarity Effects

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Abstract

Numerical solutions are presented for the growth of spherical precipitates incorporating the effect of interface curvature on local equilibrium at the interface in two-phase binary systems. This corrects an analytical solution [1].

Introduction

In recent work [1], we provided an analytical solution for the diffusion—controlled growth of a sphere including capillarity. It has been anonymously pointed out to us that this contains an inconsistency.

The concentration field in the matrix surrounding a spherical particle has to meet Fick's second law, which in spherical coordinates is:

$$\frac{\partial c}{\partial t} = \frac{D}{r^2} \frac{\partial}{\partial t} \left\{ r^2 \frac{\partial c}{\partial r} \right\} \tag{1}$$

where c is the concentration around the particle, D is the diffusion coefficient, assumed to be independent of concentration, t stands for time and r is the radial coordinate.

We proposed a solution to equation (1) as follows:

$$c\{t,\rho\} = \overline{c} + \left[\left(c^{\beta\gamma} + \frac{2c^{\beta\gamma}\Gamma}{\rho} \right) - \overline{c} \right] \frac{\phi\{r/\sqrt{Dt}\}}{\phi\{\alpha\}}$$
 (2)

where

$$\Gamma = \left(\frac{\sigma v^{\gamma}}{kT}\right) \left(\frac{1 - c^{\beta \gamma}}{c^{\gamma \beta} - c^{\beta \gamma}}\right) \tag{3}$$

and

$$\phi\{\alpha\} = \frac{1}{\alpha} \exp\left\{-\frac{\alpha^2}{4}\right\} - \frac{\sqrt{\pi}}{2} \operatorname{erfc}\left\{\frac{\alpha}{2}\right\}$$
 (4)

where \overline{c} is the average concentration of solute in the alloy, ρ is the precipitate radius, $c^{\beta\gamma}$ is the concentration of solute in the matrix (β) in equilibrium with the precipitate (γ) , σ is the surface energy per unit area, v^{γ} is the volume per atom in the precipitate phase γ , k is the Boltzmann constant, T is the temperature and $c^{\gamma\beta}$ the solute concentration of the precipitate

 (γ) in equilibrium with β , erfc is the complementary error function and $\alpha = \rho/\sqrt{Dt}$ is a growth parameter. Γ is commonly referred as the capillarity constant [2].

Equation (2) is obtained using a similarity transformation [3] but this assumes that $\phi\{\alpha\}$ is constant, which is not. We do not know how to analytically solve this difficulty, but the problem requires a solution given the need to predict the kinetics of precipitation in steels. We therefore present a numerical solution.

Method

Tanzilli and Heckel [4] have presented a numerical solution for sphere growth in the absence of capillarity. Thus, equation (1) can be expressed as

$$\frac{c_n^{j+1} - c_n^j}{\Delta t} = \frac{N - n}{L - \rho} \times \frac{c_{n+1}^j - c_{n-1}^j}{2} \times g^{j+1}
+ D \times \frac{c_{n+1}^j - 2c_n^j + c_{n-1}^j}{(L - \rho)^2/N^2} + \frac{D}{\rho + \frac{(n)(L - \rho)}{N}} \times \frac{c_{n+1}^j - c_{n-1}^j}{(L - \rho)/N} \tag{5}$$

where n=0,1,2,...,N are the nodes that divide the matrix phase in N elements each of length Δr (Fig. 1), j is a time interval, c_n^j is the concentration in n at the time interval j, Δt is the increment in time, L is the zero mass transfer boundary at the matrix, i.e. where $c_N=c_{N+1}$, and g^{j+1} the interface velocity at the time interval j+1. The time increment was set to satisfy the restriction for stable and non-oscillatory solutions to be

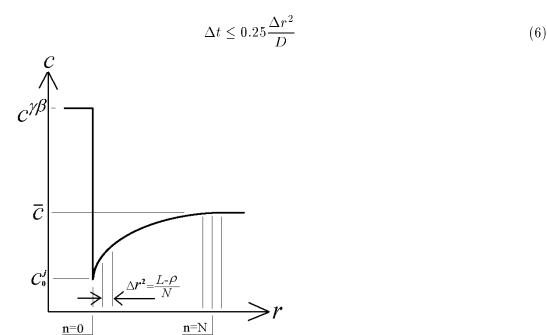


Fig. 1 Definition of finite-difference terminology.

 $r=\rho$

When capillarity effects are considered, the mass transfer equation (5) is similarly expressed as

$$\frac{\rho^{j+1} - \rho^j}{\Delta t} = \frac{D}{c^{\gamma\beta} - c_0^j} \times \frac{-c_2^j + 4c_1^j - 3c_0^j}{2(L - \rho)/N}$$
 (7)

where ρ^{j} is the particle radius at time interval j and

$$c_0^j = \overline{c} - (\overline{c} - c^{\beta \gamma}) \left(1 - \frac{\rho_c}{\rho^j} \right) \tag{8}$$

where

$$\rho_c = \frac{2c^{\beta\gamma}\Gamma}{\overline{c} - c^{\beta\gamma}} \tag{9}$$

is the critical radius, c_0^j the solute concentration at the matrix interface (Fig. 1), which is equivalent to the boundary conditions [1] met by setting the initial concentrations of all the nodes equal to \overline{c} at t=0 except c_0 , which is calculated assuming an initial particle radius of $\rho/\rho_c=1.01$. Equations (5,7,8) were thus simultaneously solved and their results are shown in Fig. 2, where the variation of the growth parameter α is plotted as a function of ρ/ρ_c for a variety of compositions (Fig. 2a); α was scaled with a starting radius of $\rho^0=1.01$ and Ω with values of $c^{\gamma\beta}=1$ and $c^{\beta\gamma}=0$. The variation of the interface velocity g with ρ/ρ_c is shown in Fig. 2b. The convergence of equations (5,7,8) was achieved when as Δt was decreased to a convenient value, a negligible change in g^j was produced, and the value of N was such that $c_N \simeq \overline{c}$.

In Fig. 2a α approaches asymptotically the value predicted by Zener's theory; this is expected as for large ρ/ρ_c the capillarity effect becomes less important. Consistent with this, the velocity (Fig. 2b) approaches a value given by $g = D\alpha^2/(2\rho)$ at large radii, while it approaches zero for small values as the driving force for growth vanishes due to capillarity.

The accuracy of the predictions given by the analytical solution [1] is illustrated in Fig. 2b where g is plotted with dotted lines against ρ/ρ_c for the indicated values of Ω . The approximate values given by the analytical method adequately predict the velocity trends; in the range of solutions observed, the maximum error was of the order of 11%; thus the analytical solution may be used for calculations where large precision is not required.

Summary

A numerical solution for the growth of particles when capillarity effects are prominent is presented. The behaviour of the solution is similar to that of an earlier (incorrect) analytical solution [1], and its precision may be acceptable for many kinetic predictions.

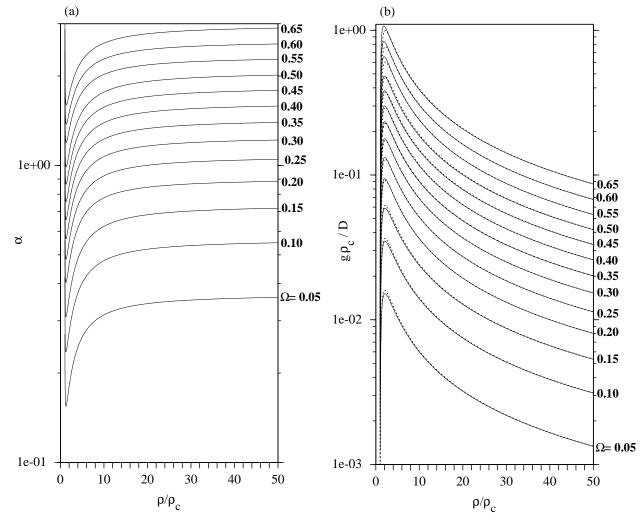


Fig. 2 Finite difference solution for (a) α and (b) g. The dotted lines represent calculations using the analytical solution.

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References

- 1. Rivera-Díaz-del-Castillo, P.E.J. and Bhadeshia, H.K.D.H.: Materials Science and Technology 17 (2001) 30–32
- Christian, J. W.: Theory of Transformations in Metals and Alloys, 2nd edn, Part I (1975)
 Oxford, Pergamon Press
- 3. Zwillinger, D.: Handbook of Differential Equations, 3rd ed, Academic Press (1998)

4. Tanzilli, R.A. and Heckel, R.W.: Transactions of the Metallurgical Society of AIME, 242 (1968) 2313–2321